



University of Minho
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Quantum Simulation of spin systems on quantum computers

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Introduction

Motivation



Figure 1: Richard Phillips Feynman (1965). *Source: The Nobel Prize*

$$N = 2^n \text{ (Complex Numbers)}$$

Beyond Classical Capacities [1], [2]

$$n = 50 \rightarrow N \approx 10^{15}$$

$$n = 300 \rightarrow N \approx 10^{90} > N_{\text{Universe}}$$

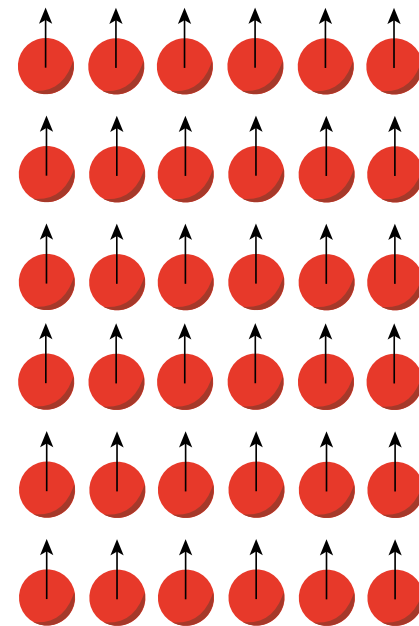


Figure 2: System of n electrons.

Quantum Simulators



Analog Quantum Simulator

$$\hat{H} \approx \hat{\mathcal{H}}$$

Digital Quantum Simulator

$$\hat{U}_{\hat{H}}(t) \approx \hat{U}_{\hat{\mathcal{H}}}(t)$$

$\hat{\mathcal{H}}$: Target Hamiltonian

\hat{H} : Simulator Hamiltonian

Designing a Quantum Simulation

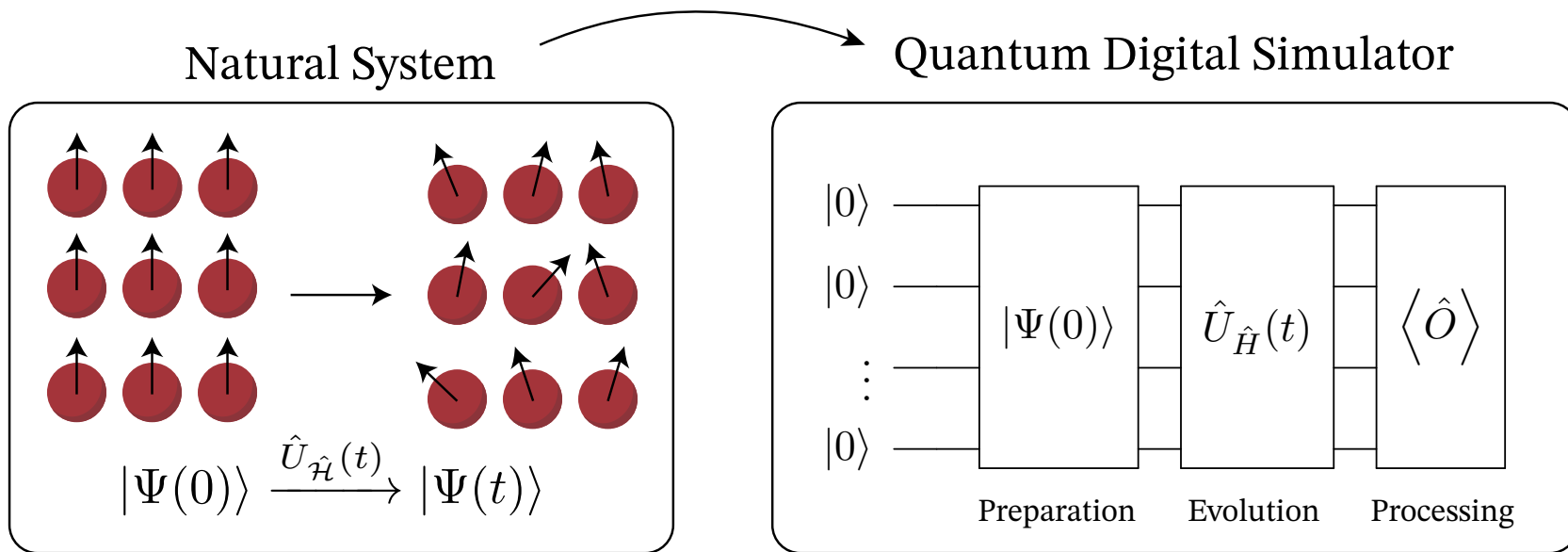
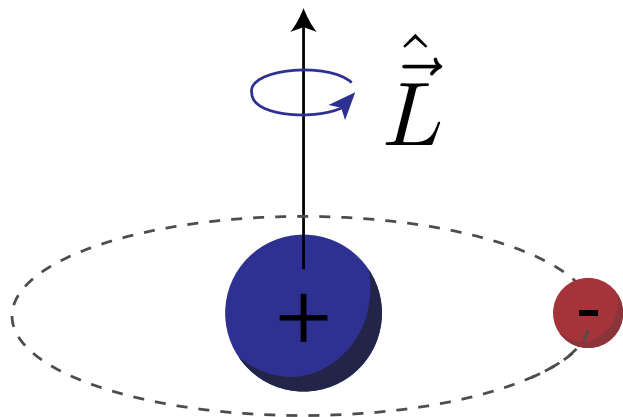
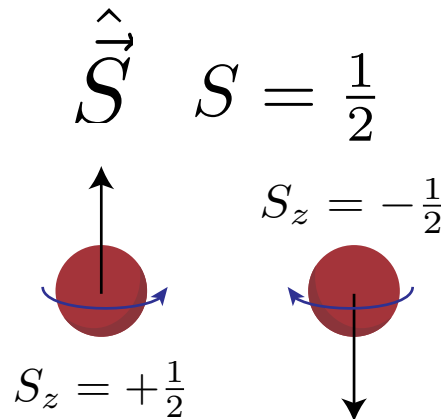


Figure 3: Scheme of digital quantum simulation using a quantum computer.

Spin Systems



Orbital Angular
Momentum



Intrinsic Angular
Momentum (Spin)

Figure 4: Diagram showing the electron's orbital angular momentum in the hydrogen atom (Left). Diagram illustrating the intrinsic angular momentum (spin) of a particle (Right) [3], [4].

Spin Chirality

Spin Chirality

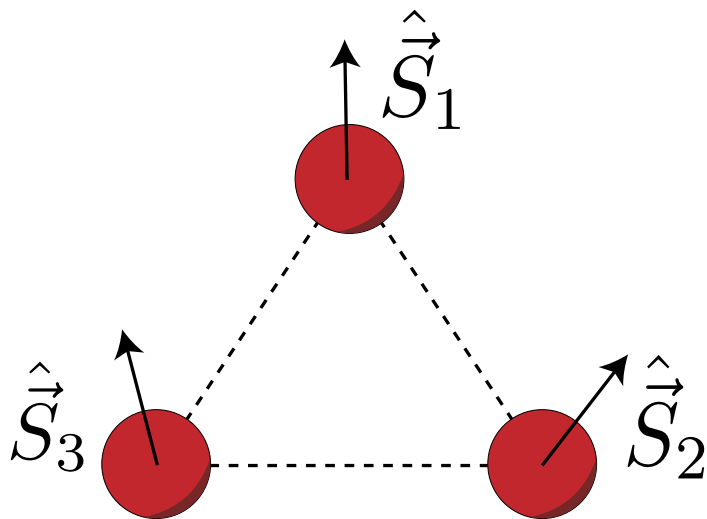


Figure 5: This figure illustrates a spin trimer and its associated spin vector operators.

Chirality Operator [5]

$$\hat{\chi} = \vec{\sigma}_1 \cdot \vec{\sigma}_2 \times \vec{\sigma}_3 \quad (1)$$

$$\vec{\sigma}_i = \left(\sigma_x^{(i)}, \sigma_y^{(i)}, \sigma_z^{(i)} \right)^T$$

Chirality eigenstates



$$\begin{aligned}
 |\chi = 0, S_z = \frac{3}{2}\rangle &= |\uparrow\uparrow\uparrow\rangle \\
 |\chi = 0, S_z = \frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|\uparrow\uparrow\downarrow\rangle + |\uparrow\downarrow\uparrow\rangle + |\downarrow\uparrow\uparrow\rangle) \\
 |\chi = 0, S_z = -\frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|\downarrow\downarrow\uparrow\rangle + |\downarrow\uparrow\downarrow\rangle + |\uparrow\downarrow\downarrow\rangle) \\
 |\chi = 0, S_z = -\frac{3}{2}\rangle &= |\downarrow\downarrow\downarrow\rangle \\
 |\chi = -2\sqrt{3}, S_z = \frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|\uparrow\uparrow\downarrow\rangle + \omega |\uparrow\downarrow\uparrow\rangle + \omega^2 |\downarrow\uparrow\uparrow\rangle) \\
 |\chi = -2\sqrt{3}, S_z = -\frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|\downarrow\downarrow\uparrow\rangle + \omega |\downarrow\uparrow\downarrow\rangle + \omega^2 |\uparrow\downarrow\downarrow\rangle) \\
 |\chi = 2\sqrt{3}, S_z = \frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|\uparrow\uparrow\downarrow\rangle + \omega^2 |\uparrow\downarrow\uparrow\rangle + \omega |\downarrow\uparrow\uparrow\rangle) \\
 |\chi = 2\sqrt{3}, S_z = -\frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|\downarrow\downarrow\uparrow\rangle + \omega^2 |\downarrow\uparrow\downarrow\rangle + \omega |\uparrow\downarrow\downarrow\rangle)
 \end{aligned} \tag{2} \quad \omega = e^{i\frac{2\pi}{3}}$$

Computational Basis



$$\boxed{|s_1 s_2 s_3\rangle \rightarrow |q_2 q_1 q_0\rangle}$$

$$\begin{aligned} |\chi = -2\sqrt{3}, S_z = \frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|100\rangle + \omega |010\rangle + \omega^2 |001\rangle) \\ |\chi = -2\sqrt{3}, S_z = -\frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|011\rangle + \omega |101\rangle + \omega^2 |110\rangle) \\ |\chi = 2\sqrt{3}, S_z = \frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|100\rangle + \omega^2 |010\rangle + \omega |001\rangle) \\ |\chi = 2\sqrt{3}, S_z = -\frac{1}{2}\rangle &= \frac{1}{\sqrt{3}} (|011\rangle + \omega^2 |101\rangle + \omega |110\rangle) \end{aligned} \tag{3}$$

$$\omega = e^{i\frac{2\pi}{3}}$$

State $|W\rangle_{N=3}$ [6], [7]



$$G(p) = \begin{pmatrix} \sqrt{p} & -\sqrt{1-p} \\ \sqrt{1-p} & \sqrt{p} \end{pmatrix} \quad (4)$$

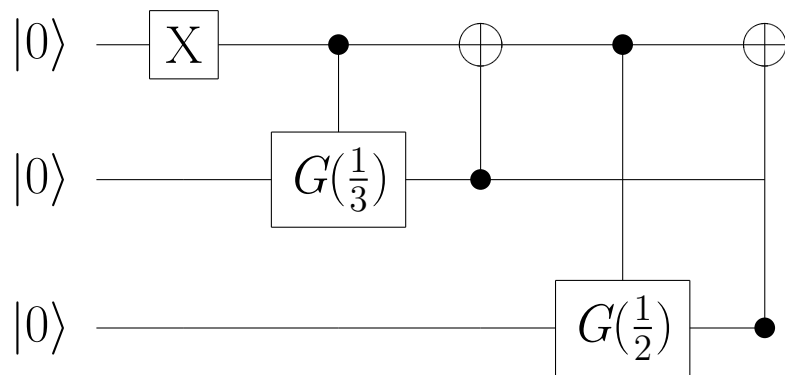


Figure 6: Quantum algorithm to prepare the $|W\rangle_{N=3}$ state.

State Preparation



$$|W\rangle_{N=3} = \frac{1}{\sqrt{N}} (|100\rangle + |010\rangle + |001\rangle)$$

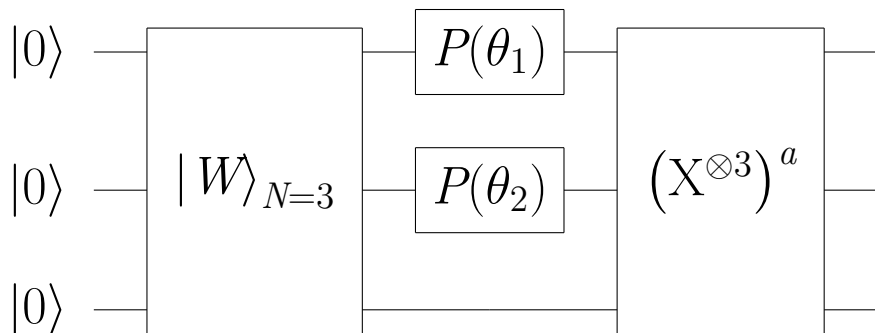


Figure 7: Quantum algorithm to prepare the non-trivial chirality operator eigenstates.

$|\chi, S_z\rangle$ Non-trivial eigenstate

Table 1: Phase gate values for each chirality eigenstate with eigenvalue χ .

χ	$P(\theta_1)$	$P(\theta_2)$
0	0	0
$-2\sqrt{3}$	ω^2	ω
$2\sqrt{3}$	ω	ω^2

$$a = \begin{cases} 1 & \text{if } S_z = -\frac{1}{2} \\ 0 & \text{if } S_z = \frac{1}{2} \end{cases} \quad (5)$$

Chirality Measurement

Chirality Operator $\hat{\chi}$ [8]



$$\begin{aligned}\hat{\chi} &= \vec{\sigma}_1 \cdot \vec{\sigma}_2 \times \vec{\sigma}_3 \\ &= \sigma_x^1 (\sigma_y^2 \sigma_z^3 - \sigma_z^2 \sigma_y^3) \\ &\quad + \sigma_y^1 (-\sigma_x^2 \sigma_z^3 + \sigma_z^2 \sigma_x^3) \\ &\quad + \sigma_z^1 (\sigma_x^2 \sigma_y^3 - \sigma_y^2 \sigma_x^3)\end{aligned}$$

$$\begin{aligned}\hat{\chi} &= 2i \left[\hat{P}_i, \hat{P}_{i+1} \right] \\ &= 2i \left(\hat{P}_i \hat{P}_{i+1} - \hat{P}_{i+1} \hat{P}_i \right)\end{aligned}$$

$$(6) \quad \hat{P}_i = 2^{\lfloor \frac{N-1}{2} \rfloor} \prod_{j=1}^{\lfloor \frac{N-1}{2} \rfloor} \left(\vec{S}_{i+j} \cdot \vec{S}_{N+i-j} + \frac{1}{4} \right)$$

$$(7) \quad \hat{P}_i = \prod_{k=0}^{\lfloor \frac{N-1}{2} \rfloor - 1} SWAP(i+k, i+N-k-2)$$

Operator $\hat{P}_i \hat{P}_{i+1}$



Let's calculate the operators for $i = 1$ and $N = 3$.

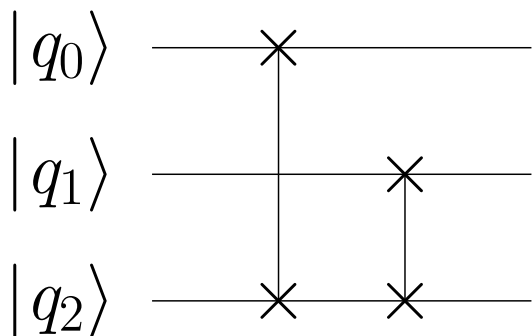


Figure 8: Operator $\hat{P}_1 \hat{P}_2$

$$\hat{P}_1 = SWAP(1, 2) \quad (8)$$

$$\hat{P}_2 = SWAP(2, 0) \quad (9)$$

$$\boxed{\hat{P}_1 \hat{P}_2 |q_2 \ q_1 \ q_0\rangle = |q_1 \ q_0 \ q_2\rangle} \quad (10)$$

Linear Combination of Unitaries [9]

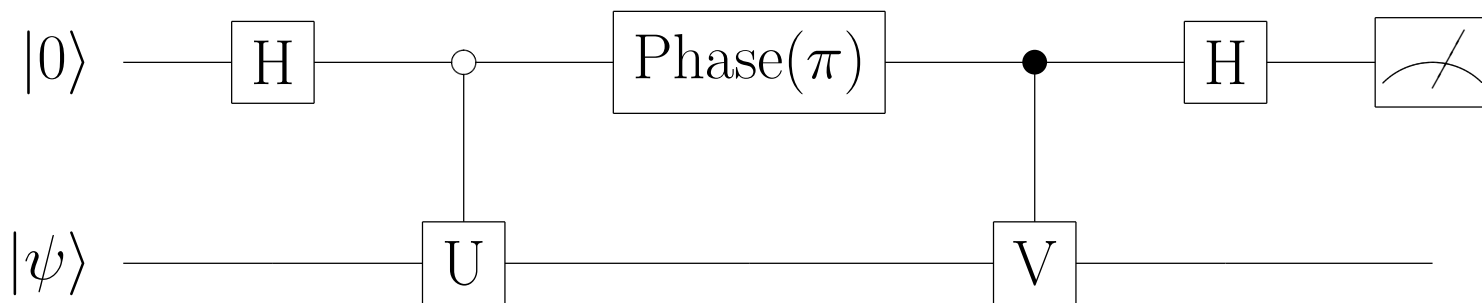


Figure 9: Linear Combination of Unitaries (LCU).

$$\text{LCU}(U, V, \pi) |0\psi\rangle = \frac{1}{2} \left[|0\rangle \otimes (U - V) |\psi\rangle + |1\rangle \otimes (U + V) |\psi\rangle \right] \quad (11)$$

$\hat{\chi}$ using LCU



$$\begin{aligned} \text{LCU}(R, R^\dagger, \pi) |0\psi\rangle &= \frac{1}{2} \left[|0\rangle \otimes (R - R^\dagger) |\psi\rangle + |1\rangle \otimes (R + R^\dagger) |\psi\rangle \right] \\ &= \frac{1}{2} \left[|0\rangle \otimes \frac{\hat{\chi}}{2i} |\psi\rangle + |1\rangle \otimes (R + R^\dagger) |\psi\rangle \right] \end{aligned} \quad (12)$$

$$\begin{aligned} U &= R = P_1 P_2 \\ V &= R^\dagger = P_2 P_1 \end{aligned} \quad (13)$$

$$\hat{\chi} = 2i(R - R^\dagger) \quad (14)$$

$$\text{Probability}(|0\rangle) = \frac{1}{16} \langle \psi | \hat{\chi}^2 | \psi \rangle \quad (15)$$

Results: χ^2



Table 2: The table summarizes the outcomes of 1000 simulations of the LCU method used to measure the value of χ^2 . Each simulation is performed with 2^{13} shots.

State	Theoretical χ^2	Simulated
$ \chi = -2\sqrt{3}\rangle$	12	12.00 ± 0.08
$ \chi = 0\rangle$	0	0.0 ± 0.0
$ \chi = 2\sqrt{3}\rangle$	12	12.00 ± 0.08

Indirect Measurement [10]

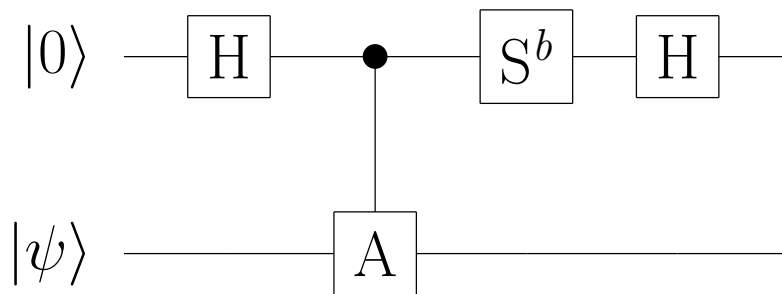


Figure 10: Simplest Hadamard test.

$$\langle Z \rangle = \begin{cases} \operatorname{Re}(\langle \psi | A | \psi \rangle) , & b = 0 \\ \operatorname{Im}(\langle \psi | A | \psi \rangle) , & b = 1 \end{cases} \quad (16)$$

Hadamard Test and LCU method

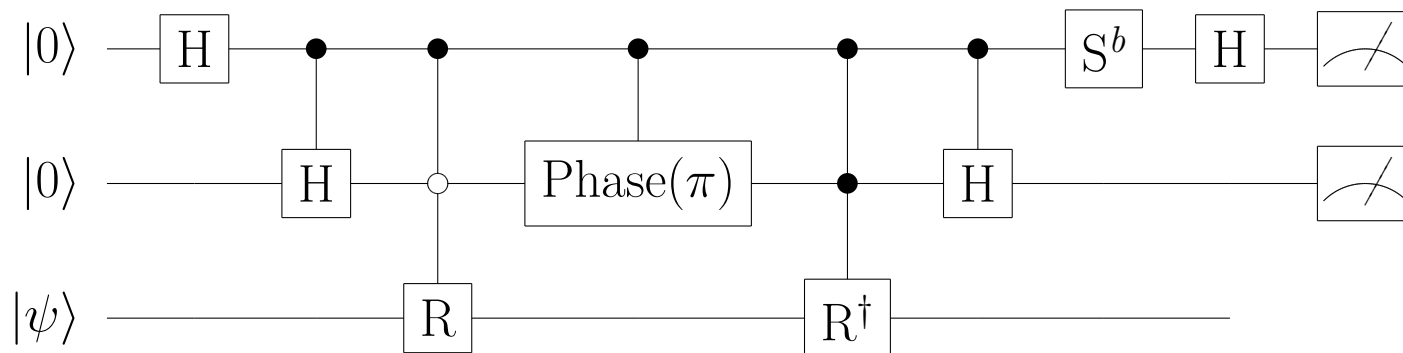


Figure 11: Hadamard test and Linear Combination of Unitaries (LCU) method to measure $\langle \hat{\chi} \rangle$.

$$p_{00} - p_{10} = \begin{cases} \frac{1}{4} \langle \psi | \text{Im} \{ \hat{\chi} \} | \psi \rangle, & b = 0 \\ \frac{1}{4} \langle \psi | \text{Re} \{ \hat{\chi} \} | \psi \rangle, & b = 1 \end{cases} \quad (17)$$

Results: $\text{Re} \{ \hat{\chi} \}$



Table 3: The table summarizes the outcomes of 1000 simulations to measure the value of χ . Each simulation is performed with 2^{13} shots.

State	Theoretical χ	Simulated
$ \chi = -2\sqrt{3}\rangle$	$-2\sqrt{3} \approx -3.46$	-3.46 ± 0.02
$ \chi = 0\rangle$	0	0.00 ± 0.03
$ \chi = 2\sqrt{3}\rangle$	$2\sqrt{3} \approx 3.46$	3.46 ± 0.02

Quantum Phase Estimation (QPE)

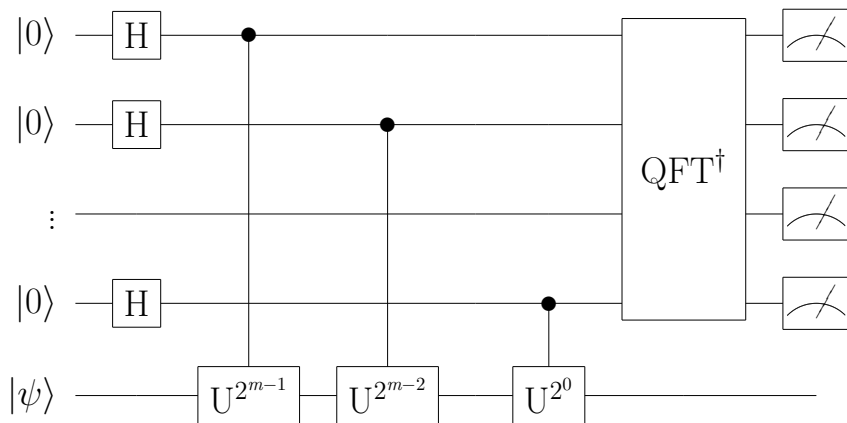


Figure 12: Quantum Phase Estimation.

$$U|\psi\rangle = e^{2\pi i\theta} |\psi\rangle \quad (18)$$

$$QPE(U)|\psi\rangle \rightarrow |2^m\theta\rangle \quad (19)$$

Chirality Time Evolution Operator



$$\begin{aligned}\hat{\chi} &= \vec{\sigma}_1 \cdot \vec{\sigma}_2 \times \vec{\sigma}_3 \\ &= \sigma_x^1 (\sigma_y^2 \sigma_z^3 - \sigma_z^2 \sigma_y^3) + \sigma_y^1 (-\sigma_x^2 \sigma_z^3 + \sigma_z^2 \sigma_x^3) + \sigma_z^1 (\sigma_x^2 \sigma_y^3 - \sigma_y^2 \sigma_x^3)\end{aligned}\quad (20)$$

$$\hat{U}_\chi(t) = e^{-it\hat{\chi}} = \lim_{n \rightarrow \infty} \left(\prod_{j,k,l} \exp \left\{ -i \frac{t}{n} \epsilon_{jkl} \sigma_j^1 \sigma_k^2 \sigma_l^3 \right\} \right)^n \quad (21)$$

$$QPE(\hat{U}_\chi(t))$$



$$e^{2\pi i\theta} = e^{-it\chi}$$

$$2\pi i\theta = 2\pi - t\chi$$

$$\theta = \left(1 - \frac{t\chi}{2\pi}\right) \bmod(1) \quad (22)$$

We want to distinguish 3 eigenvalues, positive, negative and zero.

$$m \geq \log_2(3) \approx 1.58$$

$$m = 2 \quad \text{Ancillary qubits} \quad (23)$$

Applying $\hat{U}_\chi(t)$ to Basis States



The application of the time evolution operator to the computational basis state yields the following results.

$$\hat{U}_\chi(t) |100\rangle = c_0(t) |100\rangle + c_1 |010\rangle + c_{-1} |001\rangle \quad (24)$$

$$c_j(t) = \frac{1}{3} + \frac{2}{3} \cos \left(2\sqrt{3}t + j\frac{2\pi}{3} \right) \quad (25)$$

Simplifying $\hat{U}_\chi(t)$



$$\hat{U}_\chi(t) |100\rangle = c_0(t) |100\rangle + c_1(t) |010\rangle + c_{-1}(t) |001\rangle \quad (26)$$

$$c_j(t) = \frac{1}{3} + \frac{2}{3} \cos \left(2\sqrt{3}t + j\frac{2\pi}{3} \right) \quad (27)$$

If $t = T = \frac{2\pi}{3} \frac{1}{2\sqrt{3}}$, then

$$c_{-1}(T) = 1, \quad c_0(T) = 0, \quad c_1(T) = 0 \quad (28)$$

$$\hat{U}_\chi(T) |q_2 q_1 q_0\rangle = |q_1 q_0 q_2\rangle \quad (29)$$



Hadamard Test $\hat{U}_\chi(T)$

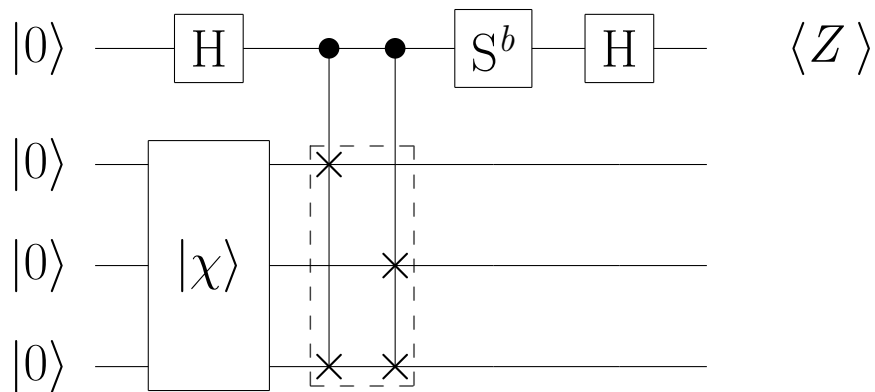


Figure 13: Hadamard test of $\hat{U}_\chi(T)$.

$$\langle Z \rangle = \begin{cases} \cos(T\chi), & b = 0 \\ \sin(T\chi), & b = 1 \end{cases} \quad (30)$$

Cycle test to measure Bargmann invariants [11].

Results: Hadamard Test $\hat{U}_\chi(T)$



Table 4: The table summarizes the outcomes of 1000 simulations to measure the value of χ . Each simulation is performed with 2^{13} shots.

State	Theoretical χ	Simulated
$ \chi = -2\sqrt{3}\rangle$	$-2\sqrt{3} \approx -3.46$	-3.46 ± 0.04
$ \chi = 0\rangle$	0	0.00 ± 0.03
$ \chi = 2\sqrt{3}\rangle$	$2\sqrt{3} \approx 3.46$	3.46 ± 0.04

$QPE(\hat{U}_\chi(T))$ 

$$\theta = \left(1 - \frac{T\chi}{2\pi}\right) \text{mod}(1), \quad T = \frac{2\pi}{3} \frac{1}{2\sqrt{3}} \quad (31)$$

$$\theta = \begin{cases} \frac{1}{3}, & \chi = -2\sqrt{3} \\ 0, & \chi = 0 \\ \frac{2}{3}, & \chi = 2\sqrt{3} \end{cases} \quad (32)$$

Results: $QPE(\hat{U}_\chi(T))$

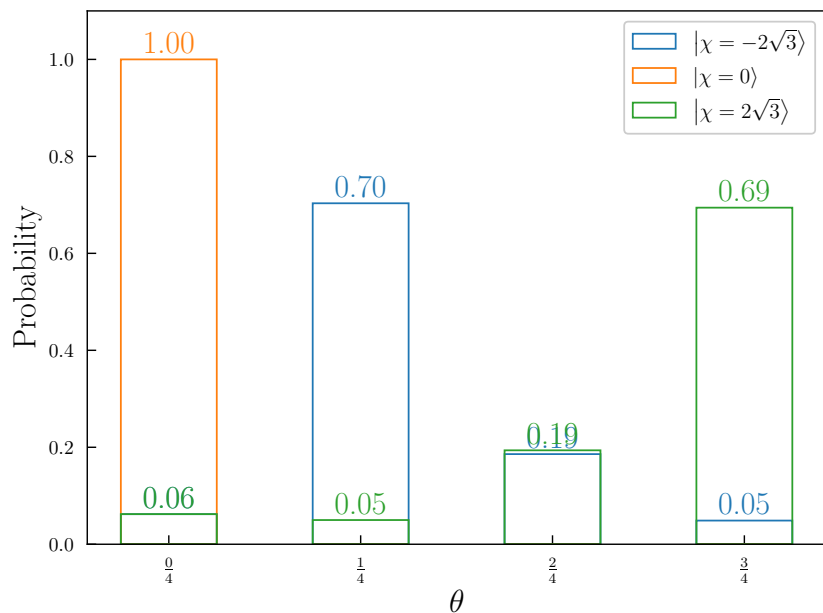


Figure 14: The graphic summarizes the outcomes of the simulation with 2^{13} shots to measure θ .

Table 5: Each state's theoretical result is presented in the table.

State	Theoretical θ	Measured
$ \chi = -2\sqrt{3}\rangle$	$\frac{1}{3} \approx 0.3\bar{3}$	0.25
$ \chi = 0\rangle$	0	0
$ \chi = 2\sqrt{3}\rangle$	$\frac{2}{3} \approx 0.6\bar{6}$	0.75

Results: Trotterized $QPE(\hat{U}_\chi(T))$

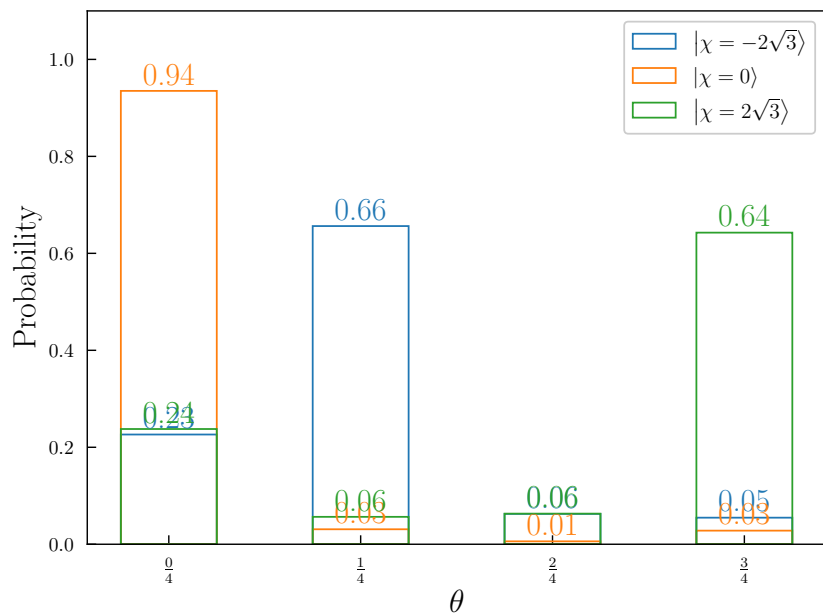


Figure 15: The graphic summarizes the outcomes of the simulation with 2^{13} shots to measure θ . This simulation used 1 trotter layer.

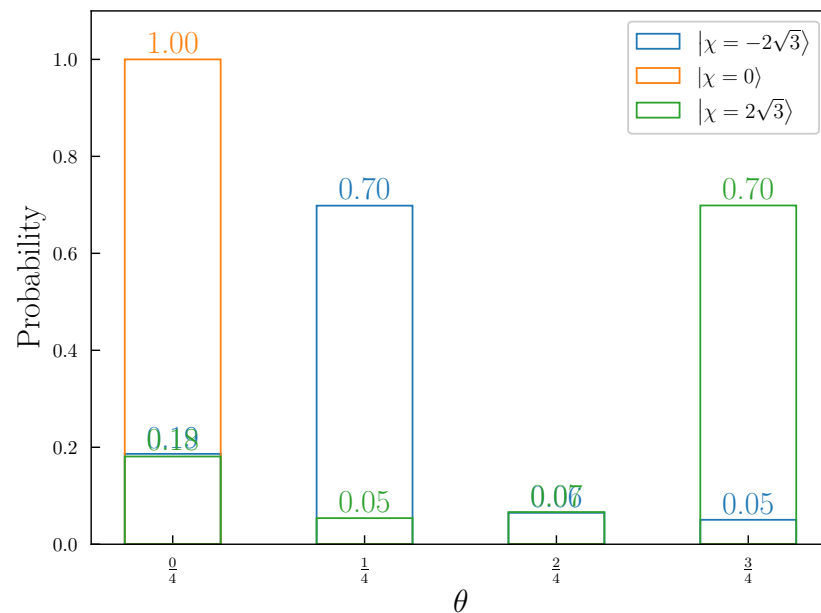


Figure 16: The graphic summarizes the outcomes of the simulation with 2^{13} shots to measure θ . This simulation used 100 trotter layers.

Conclusions

Conclusions



- ▶ Successfully implemented the chirality eigenstates.
- ▶ Successfully developed quantum algorithms for measuring the chirality property of a given state without destroying it.
- ▶ The algorithms utilize different techniques including Hadamard Test, LCU, QPE, and Trotterization.

Future Work



- ▶ Analyze and compare the algorithms.
- ▶ Implement an alternative spin system.
- ▶ Improve the algorithms with respect to gate count.
- ▶ Test the algorithms on a physical NISQ device.

Thank you for your attention



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